

"2nd Quantization" + STR

Recall: In Dirac's "2nd Quant." approach,

- $|n\rangle \rightarrow$ state with n photons.
- $a^\dagger \rightarrow$ "creates" a photon.
- $a \rightarrow$ "destroys" a photon.

→ However, the algebra is ~~ller~~ to SHO!

~~First~~

→ Now, our goal is to find a first-principles approach to finding the 2nd quantization algebra.

→ Let's start with a classical theory of SHO:

$$\Rightarrow m \frac{d^2x}{dt^2} + kx = 0$$

$$x(t) \sim e^{i\omega t}, \quad \omega \equiv \sqrt{\frac{k}{m}}$$

Constructing

the
Hamiltonian

$$H = \frac{p^2}{2m} + \frac{1}{2} m\omega^2 x^2$$

"First" ↓ w/ Quantization

$$\boxed{\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2} m\omega^2 \hat{x}^2}$$

(1)

→ Then using the canonical commutation relation,

$$[\hat{x}, \hat{p}] = i \quad [h=1, c=1]$$



$$[a, a^\dagger] = 1$$

where $a = \sqrt{\frac{m\omega}{2}} \left(\hat{x} + \frac{i\hat{p}}{m\omega} \right); a^\dagger = \sqrt{\frac{m\omega}{2}} \left(\hat{x} - \frac{i\hat{p}}{m\omega} \right)$

Q Is this consistent with STR ??

→ Check: \hat{H} transformation under ~~Lorentz~~ trans?

→ Check: The eq. of motion from \hat{H} :

$$\Rightarrow \hat{H}\psi = E\psi$$

How does this transform under Lorentz transformation?

$$\Rightarrow \cancel{\partial} \left(-\nabla^2 + \frac{m\omega^2}{2} x^2 \right) \psi = i\hbar \frac{\partial \psi}{\partial t}$$

H/W Find out how this eq. behaves when
 $x \rightarrow \gamma(x-vt); t \rightarrow \gamma(t-vx)$

Also, suppose $\psi(x, t) = e^{\pm i \vec{p} \cdot \vec{x}} e^{\pm i E t}$

$$\Rightarrow \hat{H}\psi = \cancel{i\hbar} \frac{\partial \psi}{\partial t}$$

$$\Rightarrow \left(-\nabla^2 + \frac{m\omega^2}{2} x^2 \right) \psi = i \cancel{\hbar} \frac{\partial \psi}{\partial t}$$

$$\Rightarrow \boxed{p^2 + \frac{m\omega^2}{2} x^2 = E} \quad \text{Non-relativistic energy-momentum relation!}$$

Not consistent with STR!!

Q How to construct a relativistic version of the Harmonic oscillator??

→ The simplest Lorentz-invariant eq. of motion would be: (modifying the TDSE):

[Taking hint from the behaviour of TDSE under L.T.]:

$$\cancel{\nabla^2} \boxed{\nabla^2 \psi = \frac{\partial^2}{\partial t^2} \psi}$$

where $V(x) = 0$
for convenience.

H/W Check for a 1+1D case; that

$$\frac{\partial^2}{\partial x^2} - \frac{\partial^2}{\partial t^2} = \frac{\partial^2}{\partial x'^2} - \frac{\partial^2}{\partial t'^2}$$

whereas $\star \begin{pmatrix} x' \\ t' \end{pmatrix} = \Lambda \begin{pmatrix} x \\ t \end{pmatrix}$

(3)

Therefore, in the tensor notation,

$$-\frac{\partial^2}{\partial t^2} + \frac{\partial^2}{\partial x_i^2} = \partial_\mu \partial^\mu \equiv \square$$

\downarrow
 $\eta_{\mu\nu} \partial^\mu \partial^\nu$ de'Alembertian

where

$$\partial_\mu \equiv \left(\frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$$

⇒ The ~~less~~ simplest Lorentz invariant equation of motion:

$$\square \phi(x, t) \equiv \boxed{\square \phi(x^\mu)} = 0$$

→ Repeating this calculation with a potential term, say $m^2 \phi$:

$$\boxed{(\square - m^2) \phi(x) = 0} \quad \leftarrow \text{Klein-Gordon eq.}$$

NOTE: The above equation is classical !!

⑧ So what is $\phi(x)$?

→ $\phi(x) = \phi(x^\mu) = \phi(\vec{x}, t)$ is called a "field".

→ $\phi(x)$ is a scalar quantity i.e. it is invariant under Lorentz transformation.
 $\phi(x) \rightarrow \boxed{\phi(\lambda x) = \phi(x)}$

→ $\phi(x)$ is known as the Klein-Gordon field.

→ Obvious questions:

(Q) What is the Hamiltonian/Lagrangian of $\phi(x)$?

(Q) Is $\phi(x)$ a wavefunction?

~~Is~~

→ Let's start by answering the second question:

Consider the massless KG eq:

$$\square \phi = 0 \quad \leftarrow$$

Let, $\phi(x) = a_p(t) e^{i\vec{p} \cdot \vec{x}}$

$$\Rightarrow (\partial_t^2 - \nabla^2) a_p(t) e^{i\vec{p} \cdot \vec{x}} = 0$$

$$\Rightarrow (\partial_t^2 + \vec{p}^2) a_p(t) = 0 \quad \leftarrow \begin{array}{l} \text{EoM of} \\ \text{SHO!} \\ \text{But, with } |\vec{p}|^2 \\ \text{as parameter.} \end{array}$$

∴ For a given parameter $|\vec{p}|, \omega$

KG eq. \leftrightarrow SHO

⇒ For all $|\vec{p}|$: KG eq. $\leftrightarrow \sum_{\text{all } \vec{p}} (\text{SHO})_p$

$$\Rightarrow \phi(x) = \int \frac{d^3 p}{(2\pi)^3} \left[a_p(t) e^{i\vec{p} \cdot \vec{x}} + a_p^*(t) e^{-i\vec{p} \cdot \vec{x}} \right]$$

(5)

OR

$$\boxed{\phi(x) = \int \frac{d^3 p}{(2\pi)^3} (a_p e^{-ip_\mu x^\mu} + a_p^* e^{ip_\mu x^\mu})}$$

where

$$\textcircled{a} \quad p_\mu = (\omega_{p\vec{p}}, \vec{p}) ; \quad \omega_p = |\vec{p}|$$

$$\textcircled{b} \quad p_\mu x^\mu = \omega_p x^0 - \vec{p} \cdot \vec{x}$$

⑧ Does $\phi(x)$ look like a wavefunction?

→ Not quite!

→ a_p, a_p^* ~~were~~ are upgraded to be the operators that create \hookrightarrow annihilate states.

⇒ $\boxed{\phi(x)}$ cannot be a wavefunction!

→ Now, the first question:

A field appears when STR+CM is considered.
It is convenient therefore to construct a Lagrangian first, since it is easier to write down the terms that comply with symmetries.

→ Note: Since $\phi = \phi(\vec{x}, t)$

We must write a Lagrangian density; such that:

$$\boxed{L = \int d^3x \mathcal{L}[\phi, \partial_\mu \phi]}$$

For a k-a field:

$$L = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m^2 \phi^2$$

→ Scalar!
→ ~~dim[L] = ?~~

~~(1) Variate L wrt $\phi(x)$.~~

~~(2) Find the EoM.~~

~~(1) What is the dimension of L?~~

~~(2) What is the dim[ϕ]?~~

→ Deriving the EoM of k-a field:

"Euler-Lagrange eq."

Consider, $S = \int d^4x L[\phi, \partial_\mu \phi]$

$$\begin{aligned} \Rightarrow \delta S &= \int d^4x \delta L[\phi, \partial_\mu \phi] \\ &= \int d^4x \left[\frac{\delta L}{\delta \phi} \delta \phi + \frac{\delta L}{\delta (\partial_\mu \phi)} \delta (\partial_\mu \phi) \right] \\ &= \int d^4x \left(\frac{\delta L}{\delta \phi} - \partial_\mu \left[\frac{\delta L}{\delta (\partial_\mu \phi)} \right] \right) \delta \phi \end{aligned}$$

(∴ total derivatives don't contribute)

Least action principle:

$$\delta S = 0 \quad \text{for any } \delta \phi$$

$$\Rightarrow \boxed{\frac{\partial L}{\partial \phi} - \partial_\mu \left(\frac{\partial L}{\partial (\partial_\mu \phi)} \right) = 0} \quad \text{Euler-Lagrange eq.}$$

⑦

Applying the E-L eq. for k-h L:

$$\frac{\partial \mathcal{L}}{\partial \phi} = -m^2 \phi$$

$$\frac{\partial \mathcal{L}}{\partial (\partial^\mu \phi)} = \partial^\mu \phi$$

$$\Rightarrow \boxed{\partial_\mu \partial^\mu \phi + m^2 \phi = 0} \leftarrow \text{k-h eq.}$$

NOTE: This "Hamilton's principle" for fields constitutes the Classical Field Theory.

② What is the field Hamiltonian? \rightarrow for k-h field ??
→ Iter to L:

$$H = \int d^3x \boxed{\mathcal{H}[\phi, \pi]} \rightarrow \text{Hamiltonian density.}$$

& the Legendre transform of density:

$$\mathcal{H}[\phi, \pi] = \pi[\phi, \dot{\phi}] \dot{\phi} - \mathcal{L}[\phi, \partial_\mu \phi]$$

$$\mathcal{H}[\phi, \pi[\phi, \dot{\phi}]] = \pi[\phi, \dot{\phi}] \dot{\phi} - \mathcal{L}[\phi, \partial_\mu \phi]$$

where

$$\boxed{\pi[\phi, \dot{\phi}] = \frac{\partial \mathcal{L}}{\partial \dot{\phi}}} \leftarrow \text{Conjugate momenta} \quad (\dot{\phi} \equiv \partial_t \phi)$$

⑧